## Note

## Evaluation of Electron Repulsion Integrals over Gaussian Lobe Basis Functions

The most time-consuming integral in molecular energy calculations using a gaussian lobe basis set is the electron repulsion integral

$$
\begin{equation*}
I=\left[a A_{1} b B_{1}\left|r_{12}^{-1}\right| c C_{2} d D_{2}\right] \tag{1}
\end{equation*}
$$

where, for example,

$$
\begin{equation*}
c C_{2}=\exp \left[-c\left|\mathbf{r}_{2}-\mathbf{C}\right|^{2}\right] \tag{2}
\end{equation*}
$$

is a gaussian function for electron 2 centered on the point $\mathbf{C}$. Boys [1] suggested evaluating this integral from the equation

$$
\begin{equation*}
I=f(a A, b B) f(c C, d D)(a+b+c+d)^{-1 / 2} F_{0}(t) \tag{3}
\end{equation*}
$$

where

$$
\begin{align*}
f(a A, b B) & =\left(2 \pi^{5 / 2}\right)^{1 / 2}(a+b)^{-1} \exp \left[-a b|\mathbf{A}-\mathbf{B}|^{2} /(a+b)\right],  \tag{4}\\
\mathbf{P} & =(a \mathbf{A}+b \mathbf{B}) /(a+b), \\
\mathbf{Q} & =(c \mathbf{C}+d \mathbf{D}) /(c+d), \\
t & =(a+b)(c+d)|\mathbf{P}-\mathbf{Q}|^{2} /(a+b+c+d),  \tag{5}\\
F_{0}(t) & =t^{-1 / 2} \int_{0}^{t^{1 / 2}} e^{-x^{2}} d x . \tag{6}
\end{align*}
$$

Several suggestions [2] have been made for evaluating $F_{0}(t)$. For large values of $t$,

$$
\begin{equation*}
F_{0}(t)=\frac{1}{2}(\pi / t)^{1 / 2} \tag{7}
\end{equation*}
$$

with an error of less than $e^{-t} /(2 t)$. For $t>31$, Eq. (7) has a relative error of less than $10^{-14}$. For smaller values of $t, F_{0}$ may be evaluated from

$$
\begin{equation*}
F_{0}(t)=\sum_{m=0}^{k} F_{m}\left(t_{0}\right)\left(t_{0}-t\right)^{m} / m! \tag{8}
\end{equation*}
$$

where

$$
\begin{equation*}
F_{m}(t)=t^{-(m+1 / 2)} \int_{0}^{t^{1 / 2}} x^{2 m} e^{-x^{2}} d x \tag{9}
\end{equation*}
$$

If $F_{m}\left(t_{0}\right)$ is tabulated at intervals of 0.1 , then $k=6$ (i.e., a seven-term series) is sufficient to give 13 significant figures in $F_{0}(t)$ over the range $0 \leqslant t<31$. This requires the pretabulation of 2170 coefficients $F_{m}\left(t_{0}\right) / m!$.

Shipman and Christoffersen [3] have suggested use of an eight-term Chebyshev series. They showed that this required only 32 intervals and hence only 256 coefficients to reach the same accuracy as (8). Speed in evaluation of $F_{0}(t)$ would be improved, however, by use of more intervals and fewer terms in the series. Consequently, we have examined the Chebyshev series for smaller intervals. Formally the Chebyshev series may be written on the interval $t_{i} \leqslant t \leqslant t_{i+1}$ as

$$
\begin{equation*}
F_{0}(t)=\sum_{n=0}^{\infty} a_{n}(i) T_{n}(y) \tag{10}
\end{equation*}
$$

with

$$
\begin{aligned}
y & =2\left(t-m_{i}\right) / d_{i} \\
m_{i} & =\left(t_{i}+t_{i+1}\right) / 2 \\
d_{i} & =\left(t_{i+1}-t_{i}\right) \\
-1 & \leqslant y \leqslant 1
\end{aligned}
$$

The coefficients $a_{n}(i)$ are given by

$$
\begin{equation*}
a_{n}(i)=\left(2-\delta_{n, 0}\right) \pi^{-1} \int_{-1}^{1} F_{0}(t) T_{n}(y)\left(1-y^{2}\right)^{-1 / 2} d y \tag{11}
\end{equation*}
$$

which may be evaluated by numerical integration from

$$
\begin{aligned}
a_{n}(i) & =\left(2-\delta_{n, 0}\right) N^{-1} \sum_{n=0}^{N} F_{0}(z) \cos (n \pi k / N) \\
z & =t_{i}+\frac{1}{2} d_{i}[\cos (n \pi k / N)+1]
\end{aligned}
$$

and $N=20$. The prime on the sum indicates that the first and last terms are halved.
The results of examining this series were somewhat disappointing. At intervals $d_{i}=0.1$, seven terms were still required for $0<t<2$. For $2 \leqslant t \leqslant 9$ only six terms are required, however, and for $10<t<31$ only five are needed. The basic difficulty with this approach seems to be that, for large $t, F_{0}$ behaves like $t^{-1 / 2}$ which is a difficult function to approximate.

A better alternative is to abandon (3) and avoid evaluation of $F_{0}$ altogether for large arguments. Equation (3) may be rewritten as

$$
\begin{equation*}
I=f(a A, b B) f(c C, d D)\left[(a+b)(c+d)|\mathbf{P}-\mathbf{Q}|^{2}\right]^{-1 / 2} G_{0}(t) \tag{12}
\end{equation*}
$$

where

$$
G_{0}(t)=\int_{0}^{t^{1 / 2}} e^{-x^{2}} d x=t^{1 / 2} F_{0}(t)
$$

This expression for $I$ has previously been avoided since it leads to an indeterminate form $0 / 0$ as $t$ approaches zero. For large arguments, however,

$$
G_{0}(t)=\frac{1}{2} \pi^{1 / 2}
$$

TABLE I
Number of Terms to Fit $F_{0}$ and $G_{0}$ with $d=0.1^{a}$

| $t$ | $G_{0}(t)$ <br> (Chebyshev) | $F_{0}(t)$ <br> (Chebyshev) | $F_{0}(t)$ <br> (Taylor) |
| :---: | :---: | :---: | :---: |
| 1 | $>10$ | 7 | 7 |
| 2 | 8 | 6 | 7 |
| 3 | 7 | 6 | 7 |
| 4 | 6 | 6 | 7 |
| 5 | 6 | 6 | 7 |
| 6 | 6 | 6 | 7 |
| 7 | 6 | 6 | 7 |
| 8 | 6 | 6 | 7 |
| 9 | 5 | 6 | 7 |
| 10 | 5 | 5 | 7 |
| 11 | 5 | 5 | 7 |
| 12 | 5 | 5 | 7 |
| $13-17$ | 4 | 5 | 7 |
| $18-22$ | 3 | 5 | 7 |
| $23-26$ | 2 | 5 | 7 |
| $27-31$ | 1 | 5 | 7 |

${ }^{a} d$ is the interval for which a series is accurate to $3 \times 10^{-14}$.
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TABLE II
Series and Interval Size for $F_{0}$ and $G_{0}{ }^{a}$

|  | No. of terms <br> in series | Interval <br> size | No. of coefficients <br> to be stored |
| :---: | :---: | :---: | :---: |
| $0-2$ | 5 | $2^{-6}$ | 640 |
| $2-3$ | 5 | $2^{-5}$ | 160 |
| $3-6$ | 5 | $2^{-5}$ | 480 |
| $6-8$ | 4 | $2^{-6}$ | 512 |
| $8-11$ | 4 | $2^{-5}$ | 384 |
| $11-12$ | 4 | $2^{-4}$ | 64 |
| $12-14$ | 3 | $2^{-6}$ | 384 |
| $14-16$ | 3 | $2^{-5}$ | 192 |
| $16-18$ | 3 | $2^{-4}$ | 96 |
| $18-20$ | 2 | $2^{-6}$ | 256 |
| $20-21$ | 2 | $2^{-5}$ | 64 |
| $21-22$ | 2 | $2^{-4}$ | 32 |
| $22-24$ | 2 | $2^{-5}$ | 32 |
| $24-25$ | 1 | $2^{-6}$ | 64 |
| $25-26$ | 1 | $2^{-5}$ | 32 |
| $26-30$ | 1 | $2^{-3}$ | 32 |

[^0](to 14 significant figures for $t>31$ ) so (12) can be evaluated by computing only one square root rather than the two square roots required by (3). For $t<31$, $G_{0}$ may be expanded in a Chebyshev series. For $t<7$, it was found that the $G_{0}$ series converged more slowly than the $F_{0}$ series, but for $t>7$ the $G_{0}$ series converged much more rapidly. Table I shows the number of terms required in $G_{0}$ and $F_{0}$ to reach an accuracy of $3 \times 10^{-14}$ with $d=0.1$.

Clearly a compromise procedure is suggested which uses $F_{0}$ for small values of $t$ and $G_{0}$ for large values. Our final choise of intervals is given in Table II. This choice requires storing 3424 coefficients, but no more than five terms are ever needed to evaluate $I$ to an accuracy of $2 \times 10^{-14}$. The coefficients are, of course, tabulated for the actual working form of the equation,

$$
\left\{\begin{array}{l}
F_{0}  \tag{13}\\
G_{0}
\end{array}\right\}=\left(\cdots\left(t c_{k}+c_{k-1}\right) t+\cdots+c_{1}\right) t+c_{0}
$$

The coefficients are available from the authors upon request. On the average, this algorithm is about twice as fast as one using a comparable number of coefficients based on $F_{0}$ alone. No exact timing information can be given because the average time depends on the distribution of $t$ among the different intervals.
Interval sizes which were inverse powers of 2 were used because this facilitated rapid location of the coefficients. If $[t]$ is the integer part of $t$, the first coefficient may be located by the algorithm

$$
\begin{aligned}
J & =[t], \\
K & =\left[(t-[t]) 2^{N_{J}}\right], \\
\text { location } & =M_{J}+k_{J}(K-1),
\end{aligned}
$$

where $M_{J}, N_{J}$ and $k_{J}$ (the number of terms required) are tables of integers for each $0 \leqslant J<30$. The multiplication by $2^{N_{J}}$ was accomplished by addition of $N_{J}$ to the exponent field of the floating point number $t-[t]$.

## References

1. S. F. Boys, Proc. Roy. Soc. Ser. A 200 (1950), 542.
2. I. Shavitt, in "Methods in Computational Physics" (B. Alder, S. Fernbach, and M. Rotenberg, Eds.), Vol. 2, Academic Press, New York, 1963.
3. L. L. Shipman and R. E. Christoffersen, Comp. Phys. Comm. 2 (1971), 201.

Received: September 18, 1973; revised December 11, 1973
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[^0]:    ${ }^{a} F_{0}$ is used for $0<t<3 ; G_{0}$ is used for $3 \leqslant t<30 ; F_{0}=1$ is used for $t=0 ; G_{0}=\frac{1}{2} \pi^{1 / 2}$ is used for $30 \leqslant t<\infty$.

